

ALGORITHMIC ASPECTS OF A LOCALIZED BRIDGING SCALE METHOD

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ABSTRACT

In this paper, we present a modified version of the bridging scale method (BSM), recently proposed by Liu *et al.* The major purpose of this work is to propose a consistent computational BSM algorithm to maintain of a greate advantage of the element to global assembling scheme in the finite element method. In BSM, the total displacement $\mathbf{u}(\mathbf{x})$ in a domain Ω may be decomposed into coarse scale part $\bar{\mathbf{u}}(\mathbf{X})$ and fine scale part $\check{\mathbf{u}}(\mathbf{X})$:

$$\mathbf{u}(\mathbf{X}) = \bar{\mathbf{u}}(\mathbf{X}) + \check{\mathbf{u}}(\mathbf{X}).$$

The coarse scale displacement at α -th atom in e -th element $\bar{\mathbf{u}}^e(\mathbf{X}_\alpha)$ may be interpolated by the finite element nodal displacements $\hat{\mathbf{u}}^e$:

$$\bar{\mathbf{u}}^e(\mathbf{X}_\alpha) = \mathbf{N}_\alpha^e \hat{\mathbf{u}}^e.$$

We may define the difference between the total displacement and the coase scale displacement at \mathbf{X}_α :

$$\Delta_\alpha^e = \mathbf{u}_\alpha^e - \mathbf{N}_\alpha^e \hat{\mathbf{u}}^e.$$

To obtain the optimal approximation \mathbf{u} , we may define the weighted residual functional in each finite element region Ω^e :

$$\begin{aligned} \mathcal{E}^e &= \sum_{\alpha=1}^{n\alpha_e} m_\alpha^e (\Delta_\alpha^e \cdot \Delta_\alpha^e) = \sum_{\alpha=1}^{n\alpha_e} m_\alpha^e (\mathbf{u}_\alpha^e - \mathbf{N}_\alpha^e \hat{\mathbf{u}}^e)^T (\mathbf{u}_\alpha^e - \mathbf{N}_\alpha^e \hat{\mathbf{u}}^e) \\ &= \left(\mathbf{u}^e - \mathbf{N}^e \hat{\mathbf{u}}^e \right)^T \mathbf{m}^e \left(\mathbf{u}^e - \mathbf{N}^e \hat{\mathbf{u}}^e \right). \end{aligned}$$

where $n\alpha_e$ is the number of atoms in the e -th element, \mathbf{u}^e and \mathbf{m}^e are the total displacement and mass arrays of all atoms in the element respectively, and \mathbf{N}^e is the interpolation operator from the element

nodal displacements to displacements at all atoms in the element. The minimization of the element based weighted residual functional may yields following:

$$\frac{\partial \mathcal{E}^e}{\partial \overset{\circ}{\mathbf{u}}^e} = 0 \quad \rightarrow \quad \overset{\circ}{\mathbf{m}}^e \overset{\circ}{\mathbf{u}}^e = (\mathbf{N}^e)^T \mathbf{m}^e \mathbf{u}^e, \quad \overset{\circ}{\mathbf{m}}^e = (\mathbf{N}^e)^T \mathbf{m}^e \mathbf{N}^e$$

The coarse and fine scale parts of the displacement in the element are defined as

$$\begin{aligned} \bar{\mathbf{u}}^e &= \mathbf{N}^e \overset{\circ}{\mathbf{u}}^e = \mathbf{N}^e (\overset{\circ}{\mathbf{m}}^e)^{-1} (\mathbf{N}^e)^T \mathbf{m}^e \mathbf{u}^e = \mathbf{p}^e \mathbf{u}^e \\ \check{\mathbf{u}}^e &= \mathbf{u}^e - \bar{\mathbf{u}}^e = (\mathbf{I}^e - \mathbf{p}^e) \mathbf{u}^e = \mathbf{q}^e \mathbf{u}^e \end{aligned}$$

Then the total displacement is presented as:

$$\mathbf{u}^e = \bar{\mathbf{u}}^e + \check{\mathbf{u}}^e = \mathbf{N}^e \overset{\circ}{\mathbf{u}}^e + \mathbf{u}^e - \mathbf{p}^e \mathbf{u}^e = \mathbf{N}^e \overset{\circ}{\mathbf{u}}^e + \mathbf{q}^e \mathbf{u}^e$$

To bridge the scale, the total displacement in the right hand side in the above equation can be presented by the the displacement from the molecular dynamics solution $\dot{\mathbf{u}}^e$.

$$\mathbf{u}^e = \mathbf{N}^e \overset{\circ}{\mathbf{u}}^e + \dot{\mathbf{u}}^e - \mathbf{p}^e \dot{\mathbf{u}}^e$$

Differentiating the total displacement \mathbf{u}^e with respect to time gives

$$\dot{\mathbf{u}}^e = \mathbf{N}^e \overset{\circ}{\dot{\mathbf{u}}^e} + \mathbf{q}^e \dot{\dot{\mathbf{u}}^e}$$

The Lagrangian \mathcal{L} of the bridging scale method can be written as:

$$\mathcal{L} = \sum_{e=1}^m \frac{1}{2} (\overset{\circ}{\dot{\mathbf{u}}^e})^T \overset{\circ}{\mathbf{m}}^e \overset{\circ}{\dot{\mathbf{u}}^e} + \sum_{e=1}^m \frac{1}{2} (\dot{\dot{\mathbf{u}}^e})^T \dot{\mathbf{m}}^e \dot{\dot{\mathbf{u}}^e} - \sum_{e=1}^m \mathcal{U}^e(\overset{\circ}{\mathbf{u}}^e, \dot{\dot{\mathbf{u}}^e})$$

where m is the number of finite elements, \mathcal{U} is the interatomic potential energy and the fine scale mass matrix $\dot{\mathbf{m}}^e$ is defined to be $\dot{\mathbf{m}}^e = (\mathbf{q}^e)^T \dot{\mathbf{m}}^e$. The element nodal and atomic variables can be presented by the corresponding global value:

$$\overset{\circ}{\dot{\mathbf{u}}^e} = \overset{\circ}{\mathcal{A}}^e \overset{\circ}{\dot{\mathbf{U}}}, \quad \dot{\dot{\mathbf{u}}^e} = \dot{\mathcal{A}}^e \dot{\dot{\mathbf{U}}}$$

where $\overset{\circ}{\dot{\mathbf{U}}}$ and $\dot{\dot{\mathbf{U}}}$ are the global nodal and the global atomic velocity arrays respectively, $\overset{\circ}{\mathcal{A}}^e$ and $\dot{\mathcal{A}}^e$ are operators which map global value to corresponding element value. The Lagrangian can be written in the global form as:

$$\mathcal{L} = \frac{1}{2} (\overset{\circ}{\dot{\mathbf{U}}})^T \overset{\circ}{\mathbf{M}} \overset{\circ}{\dot{\mathbf{U}}} + \frac{1}{2} (\dot{\dot{\mathbf{U}}})^T \dot{\mathbf{M}} \dot{\dot{\mathbf{U}}} - \mathcal{U}(\overset{\circ}{\mathbf{U}}, \dot{\dot{\mathbf{U}}})$$

where

$$\overset{\circ}{\mathbf{M}} = \sum_{e=1}^m (\overset{\circ}{\mathcal{A}}^e)^T \overset{\circ}{\mathbf{m}}^e \overset{\circ}{\mathcal{A}}^e, \quad \dot{\mathbf{M}} = \sum_{e=1}^m (\dot{\mathcal{A}}^e)^T \dot{\mathbf{m}}^e \dot{\mathcal{A}}^e$$

Finally the coupled bridging scale equations of motion have been derived as:

$$\dot{\mathbf{M}} \dot{\dot{\mathbf{U}}} = \dot{\mathbf{F}}(\dot{\dot{\mathbf{U}}}), \quad \overset{\circ}{\mathbf{M}} \overset{\circ}{\dot{\mathbf{U}}} = \overset{\circ}{\mathbf{F}}(\overset{\circ}{\mathbf{U}}, \dot{\dot{\mathbf{U}}})$$

REFERENCES

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