

ADHESIVE FORCES ON HELIUM IN NONTRIVIAL GEOMETRIES

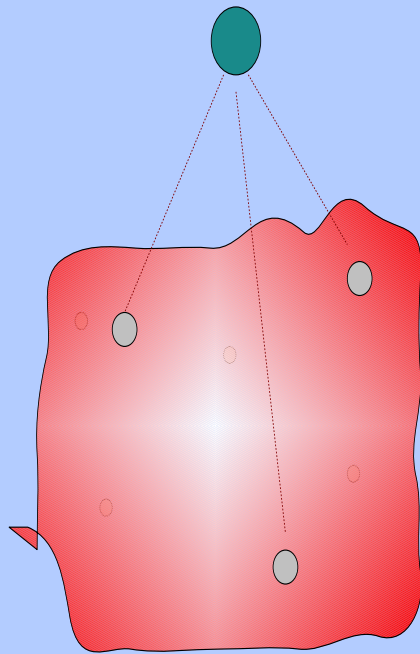
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shernand@df.uba.ar

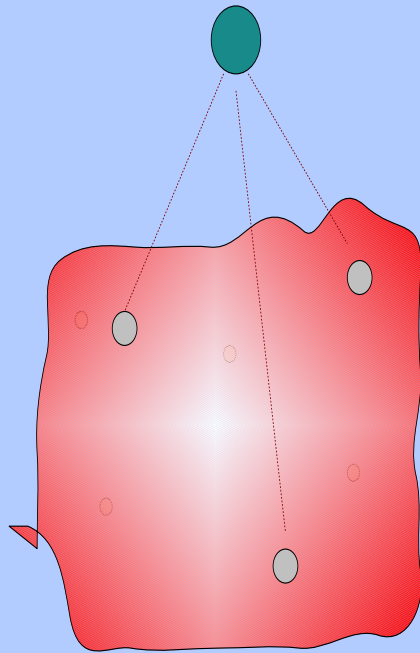
¹ *University of Buenos Aires, Argentina*

² *University of Barcelona, Spain*

Vapor particle + matter

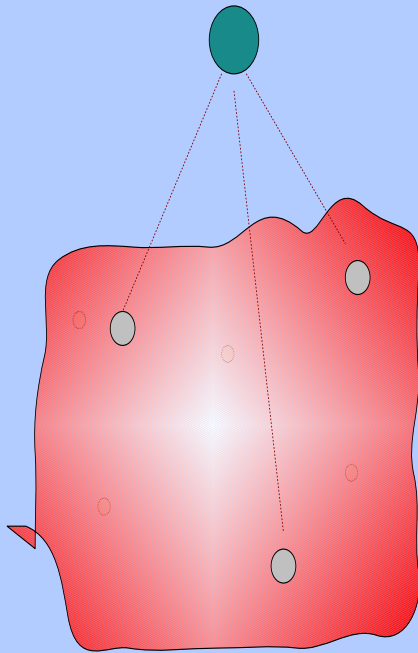


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$$U(\mathbf{r}) = \sum_i V(\mathbf{r} - \mathbf{r}')$$

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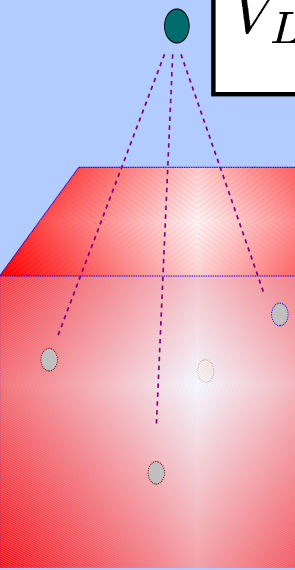
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A popular choice

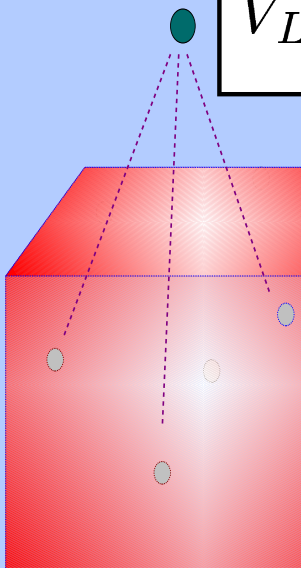
$$V_{LJ}(\mathbf{r}) \equiv V_{6-12}(\mathbf{r}) = \varepsilon \left[\left(\frac{r_{min}}{r} \right)^{12} - 2 \left(\frac{r_{min}}{r} \right)^6 \right]$$

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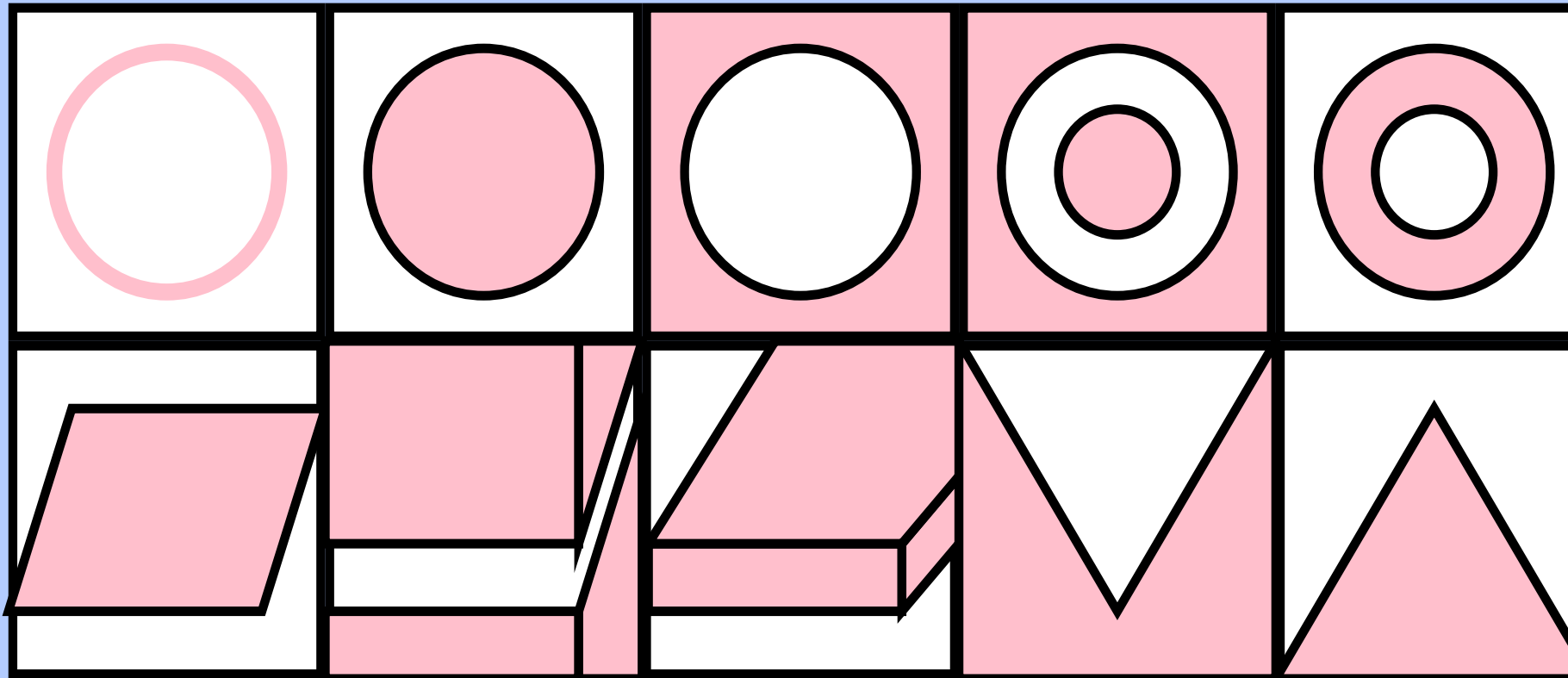
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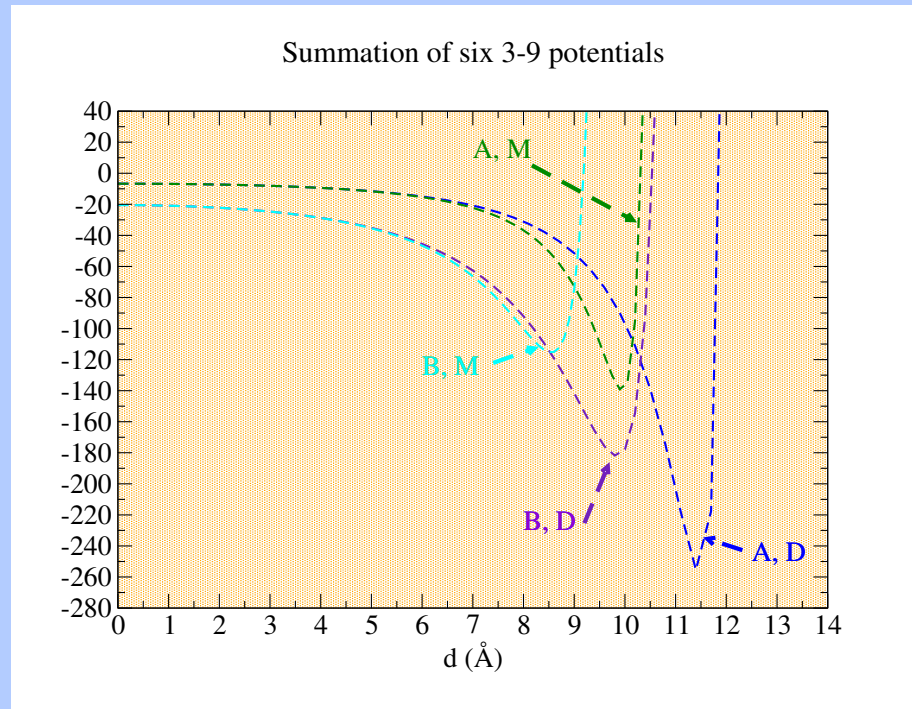
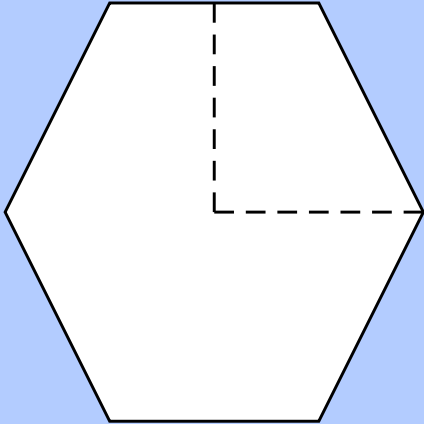
$$U(z) \equiv V_{3-9}(z) = \frac{D}{2} \left[\left(\frac{z_{min}}{z} \right)^9 - 3 \left(\frac{z_{min}}{z} \right)^3 \right]$$

Geometries with high symmetry

The LJ potential is analytically integrable (or quasiintegrable)

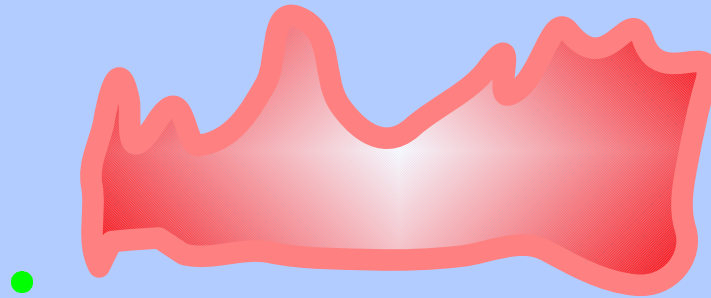


FSM-16 hexagonal pores with six 3-9's

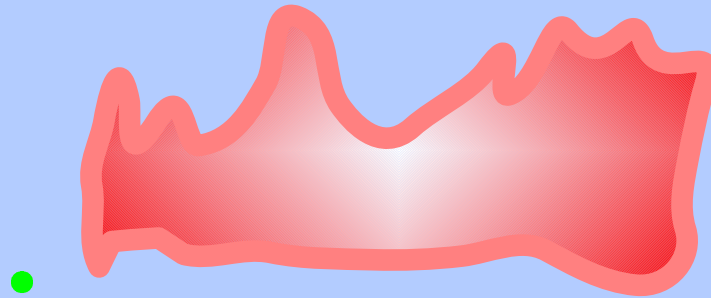


(from M. Rossi, D. E. Galli and L. Reatto, JLTP 146, 98 (2006))

Open questions

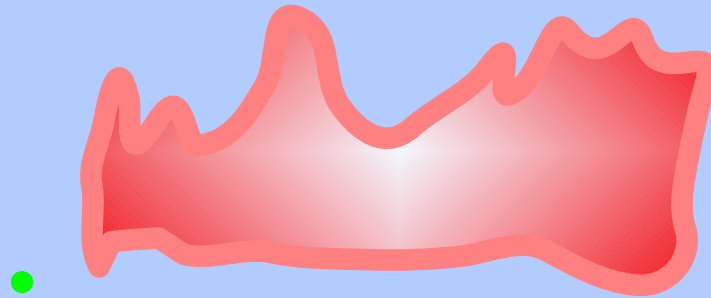


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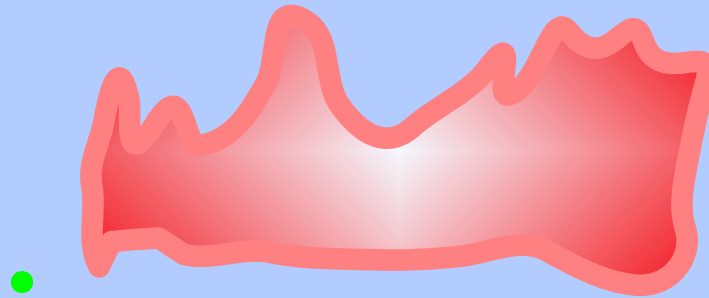
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Open questions



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- How accurate is a LJ potential?

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- How accurate is a LJ potential?
- For metallic planar half-solids, an *ab-initio* $U(z)$ is

$$U_{CCZ}(z) \neq \int d\mathbf{r}' n(\mathbf{r}') V_{pair}(\mathbf{r} - \mathbf{r}')$$

(A. Chizmeshya, M. W. Cole, and E. Zaremba, J. Low Temp. Phys. 110, 677 (1998).)

New question

Is the inverse problem solvable?, *i.e.*, given $U_{ref}(z)$, can we find sources $V(\mathbf{r} - \mathbf{r}')$?

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$$U_{ref}(z) = \int_{-\infty}^0 dz' \int \int dx' dy' V(\mathbf{r} - \mathbf{r}')$$

If so, then for *almost any* geometry,

$$U(\mathbf{r}) = \int d\mathbf{r}' n(\mathbf{r}') V(\mathbf{r} - \mathbf{r}')$$

It is solvable!

For a half-solid with bulk density ρ_0 ,

$$V(z) = \frac{U''_{ref}(z)}{2\pi\rho_0 z}$$

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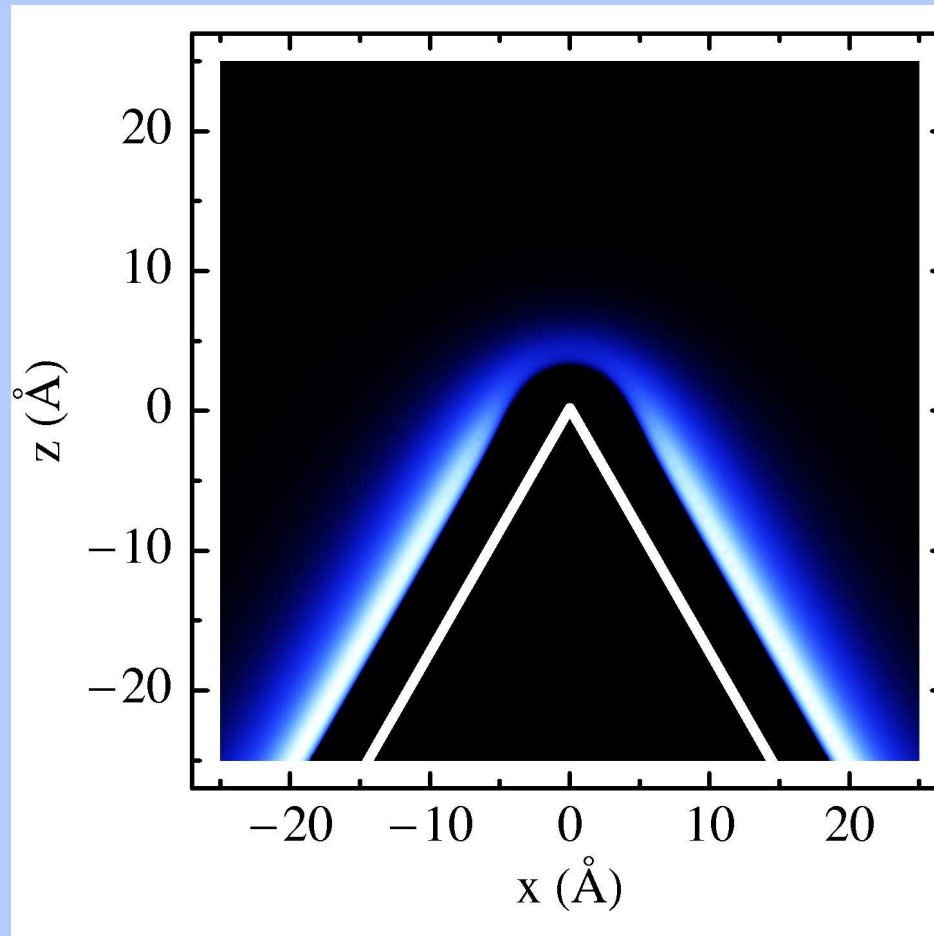
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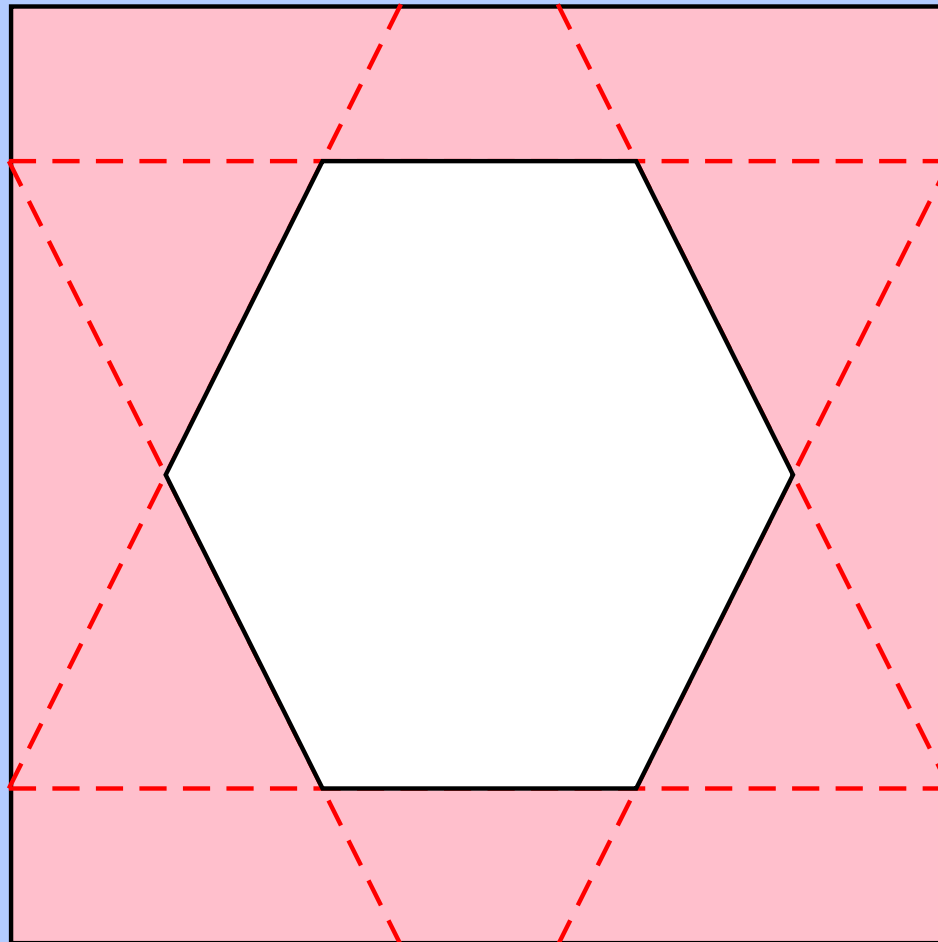
so that for other matter distributions $n(\mathbf{r}')$

$$U(\mathbf{r}) = \int d\mathbf{r}' n(\mathbf{r}') \frac{U''_{ref}(|\mathbf{r}-\mathbf{r}'|)}{2\pi\rho_0|\mathbf{r}-\mathbf{r}'|}$$

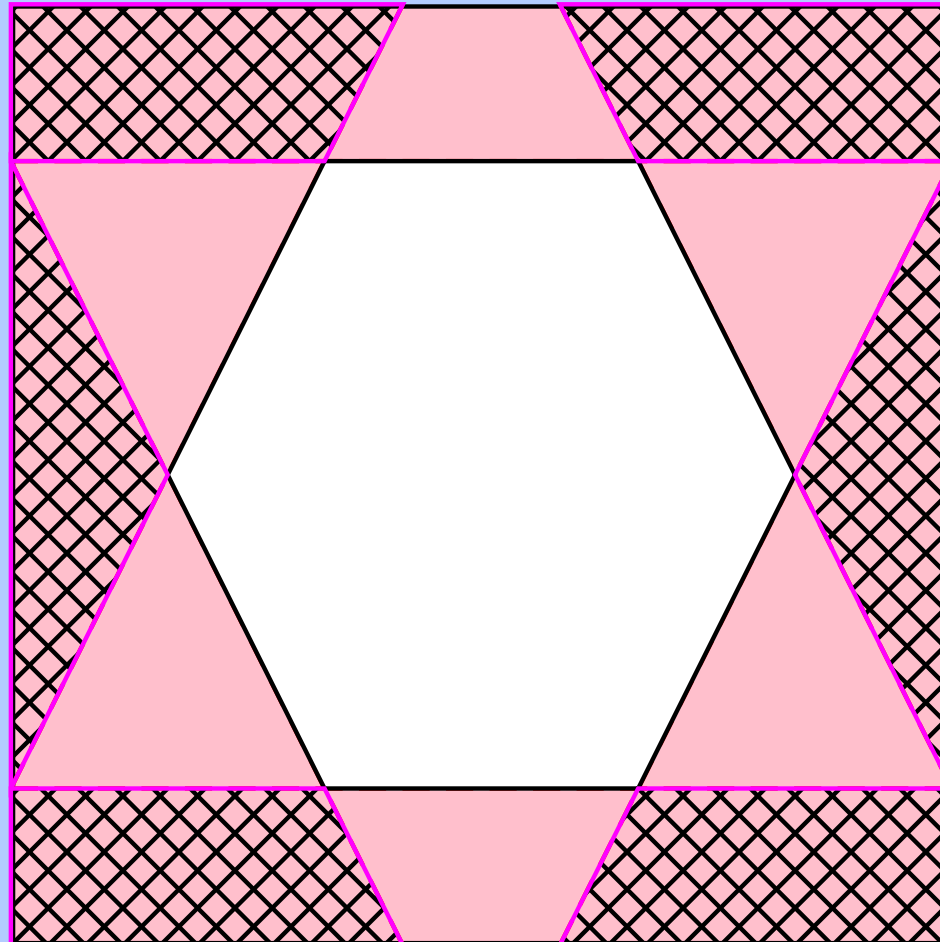
A friendly unit cell



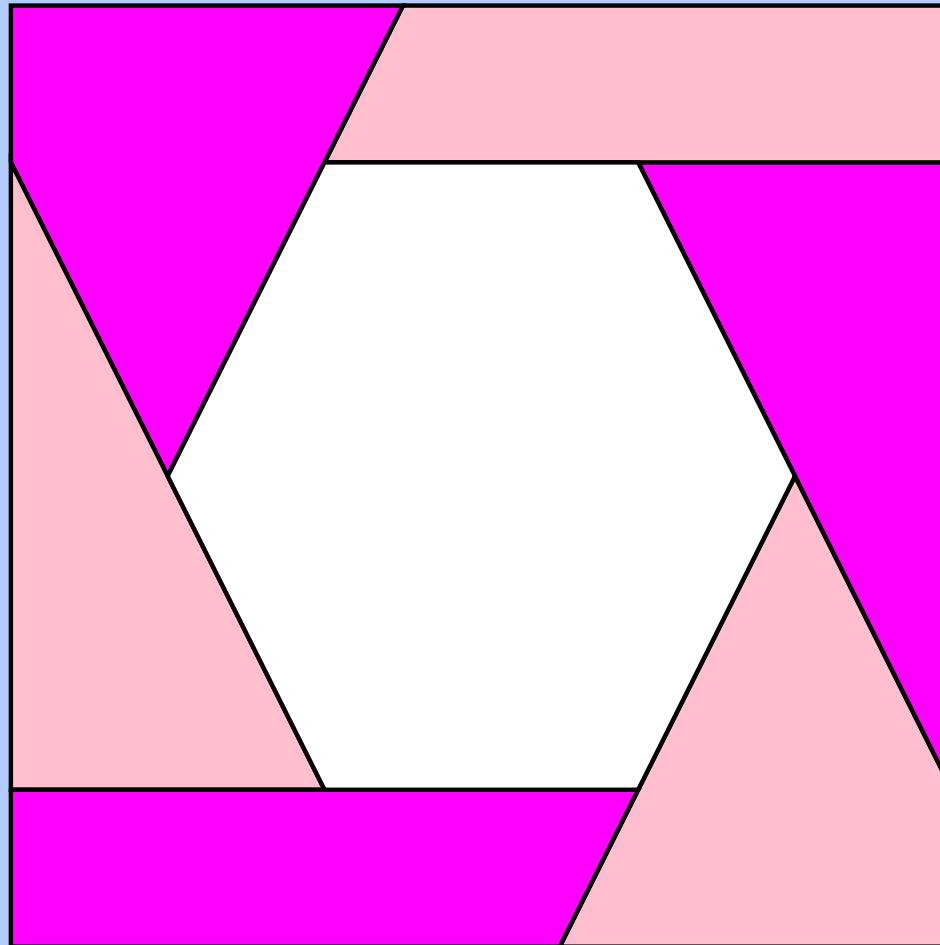
Six half-solids make an hexagonal pore



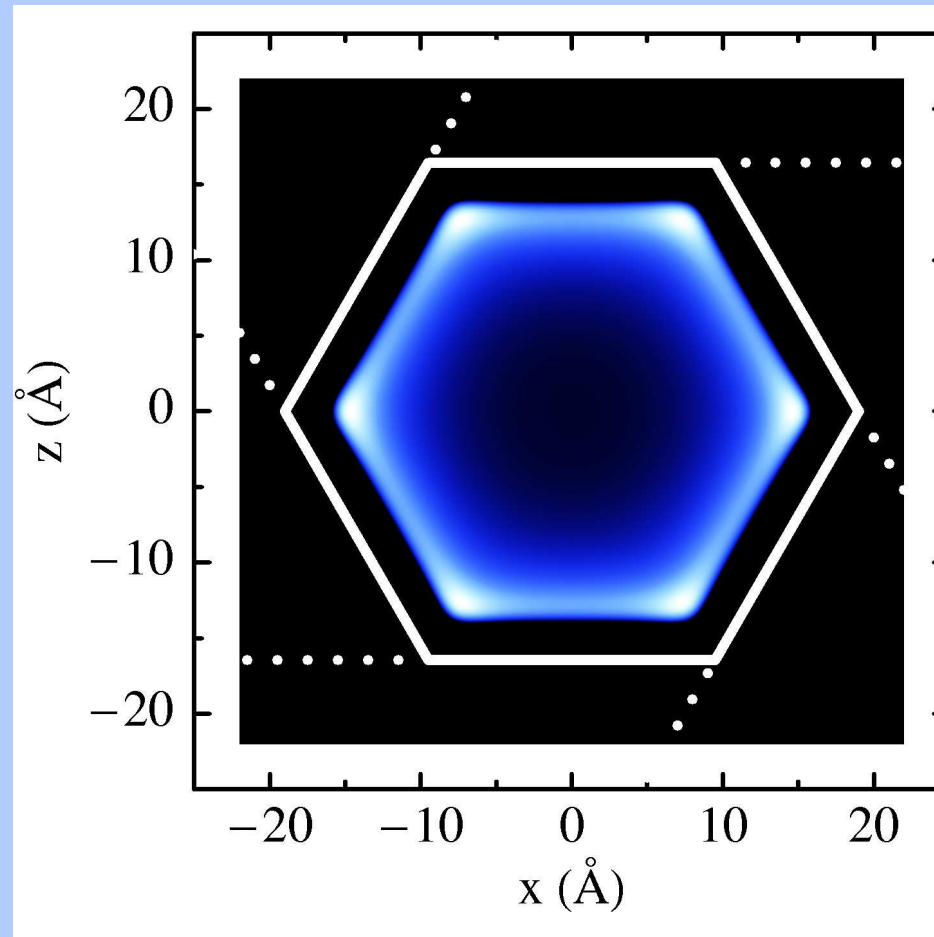
with strong overlap at the vertices



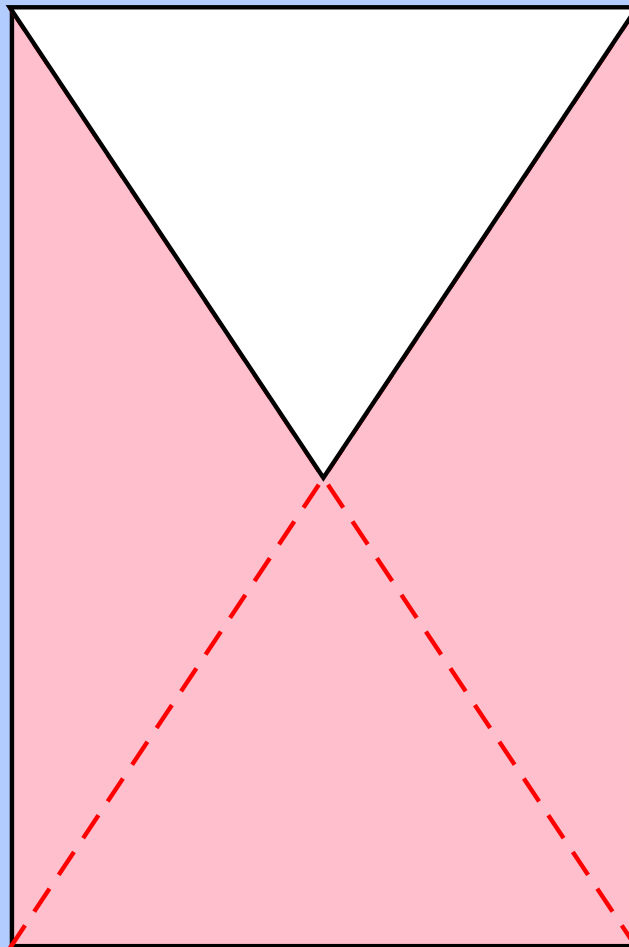
so, we recommend SIX CUSPS INSTEAD



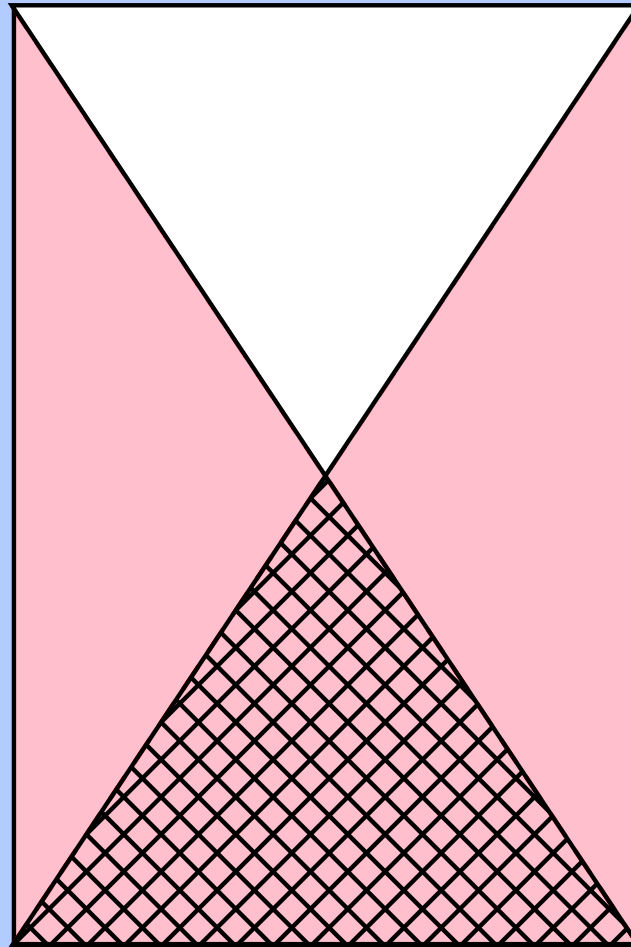
with potential landscape



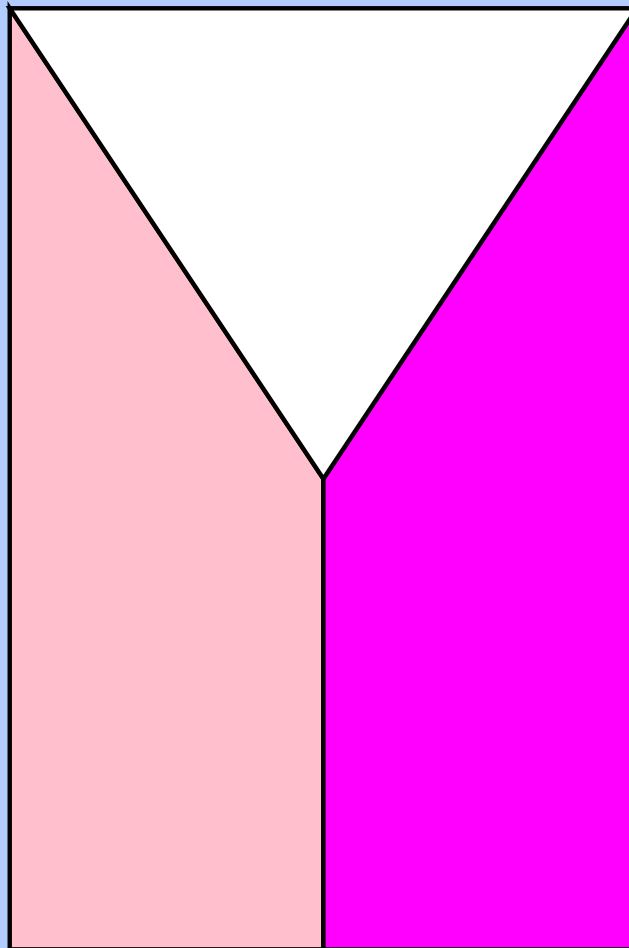
The wedge



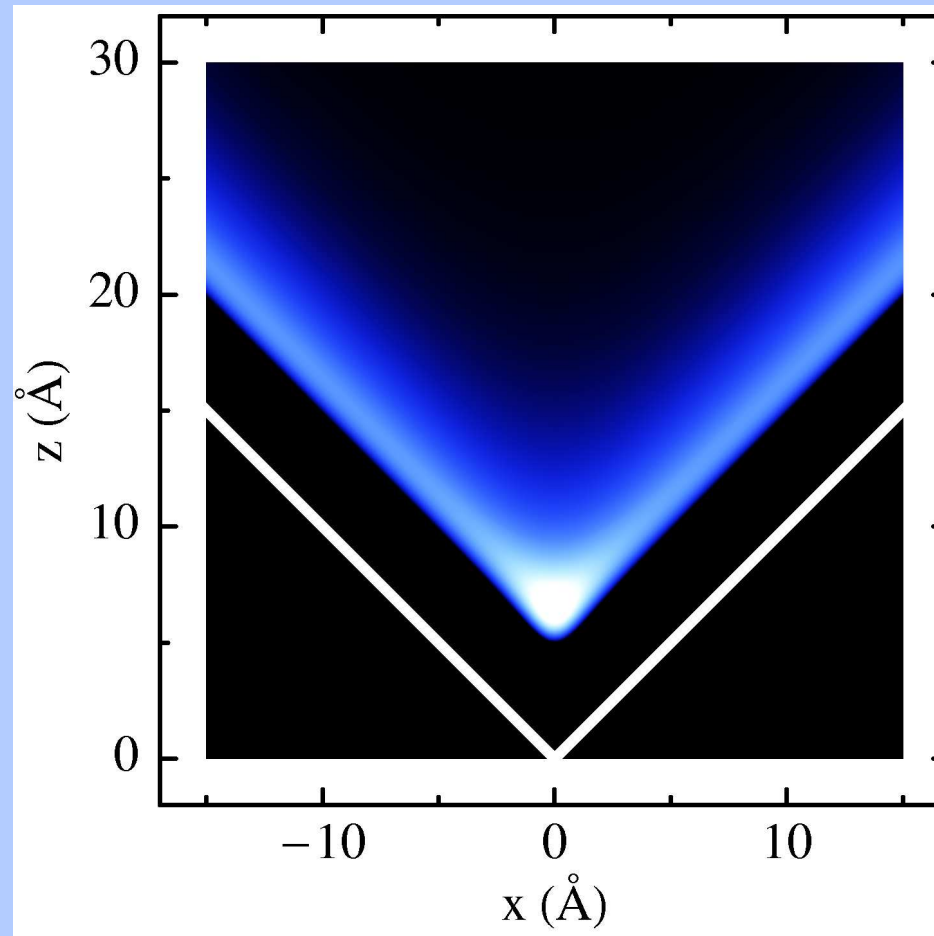
The wedge was (ESH *et al* in PRB 73, 245406 (2006))



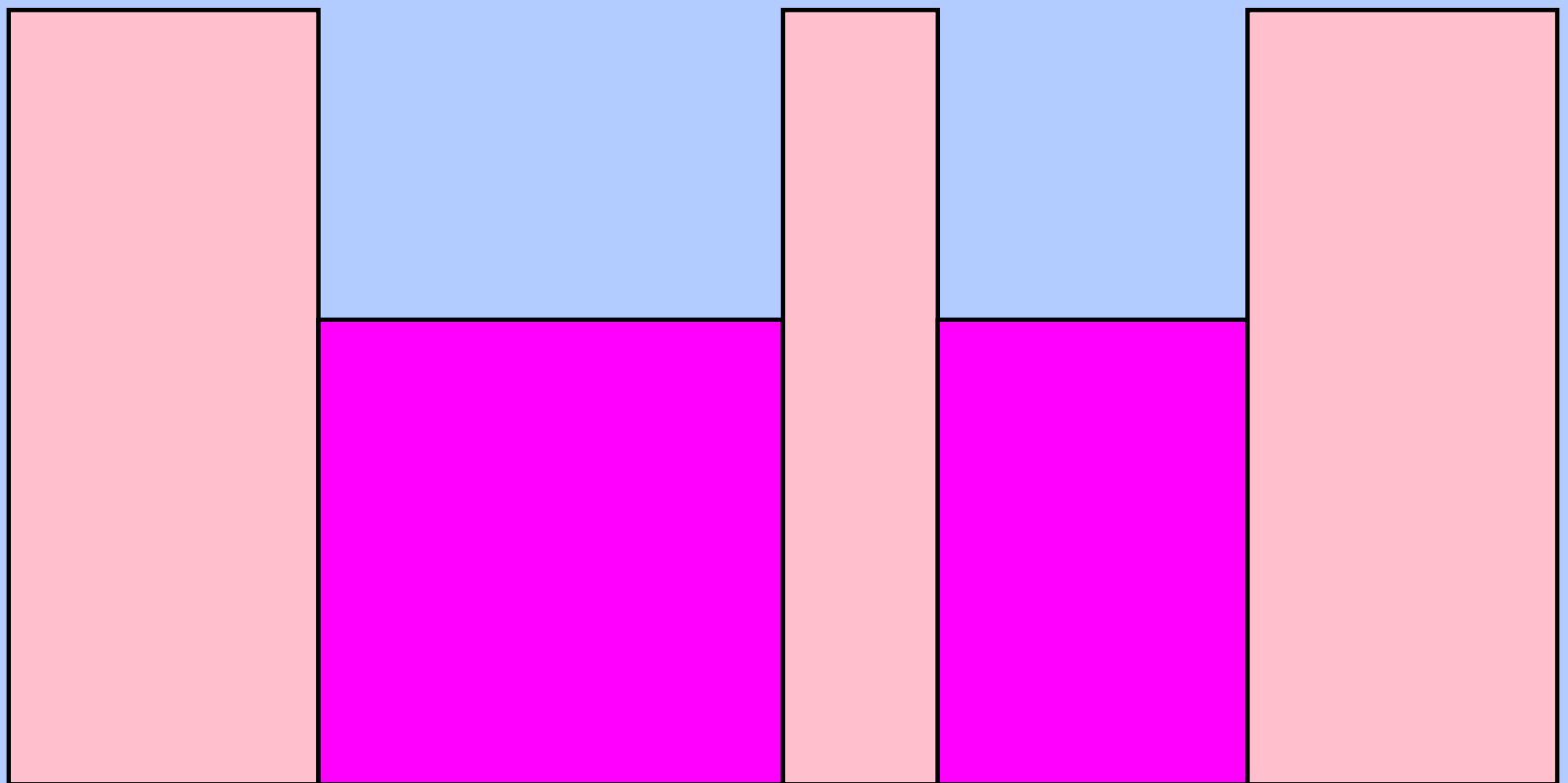
while it may be



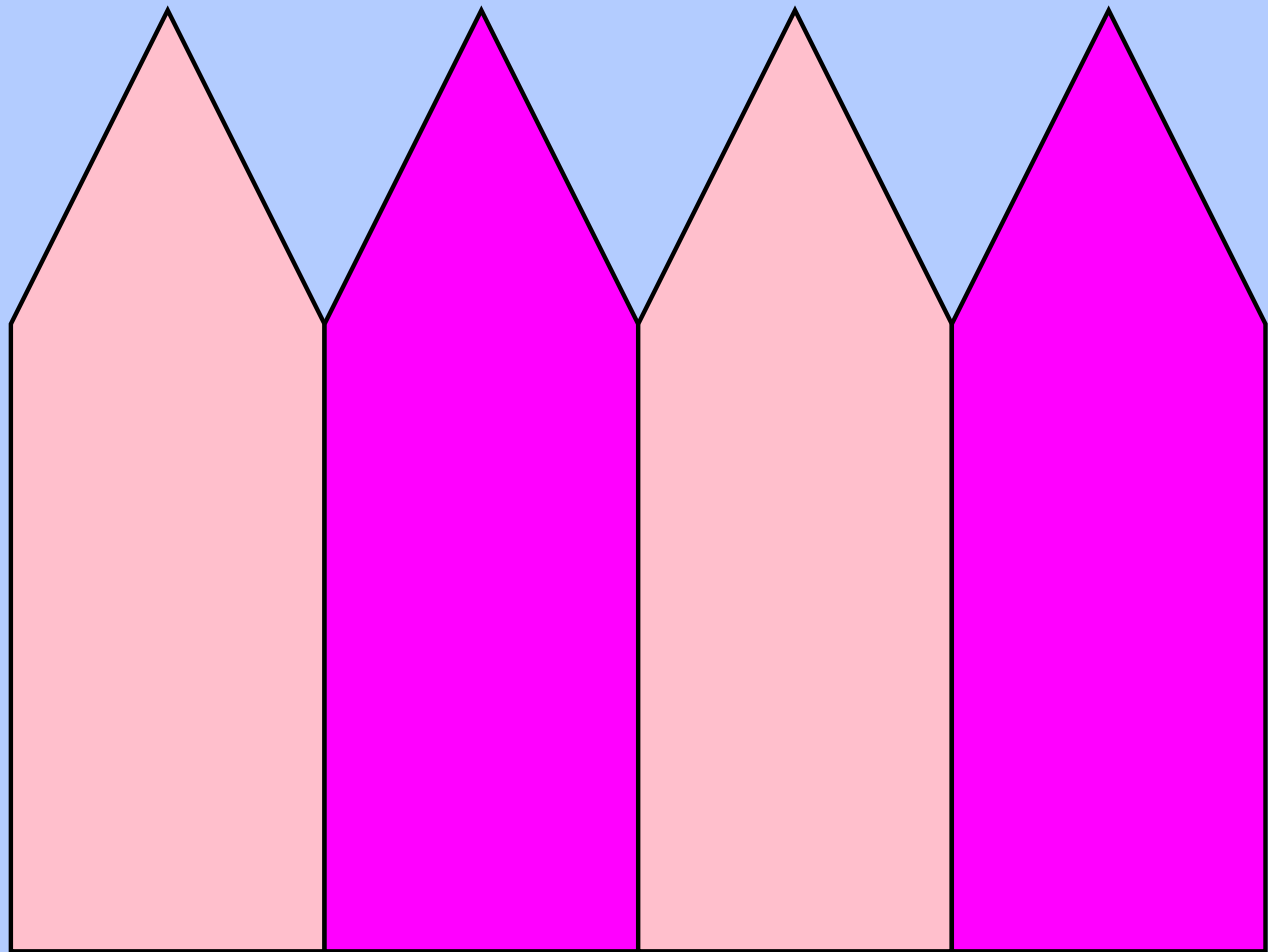
with potential



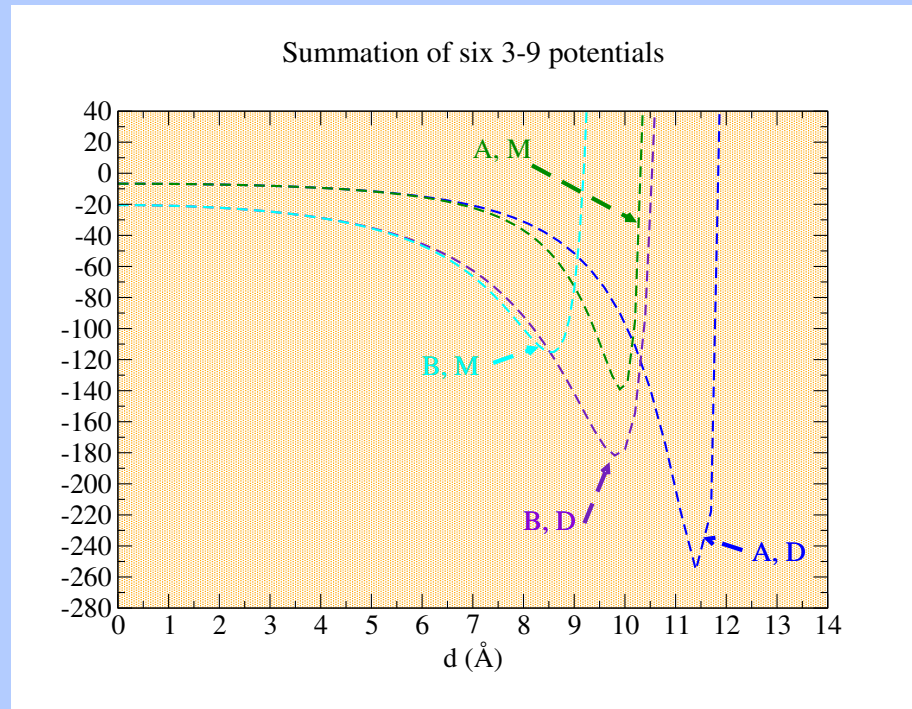
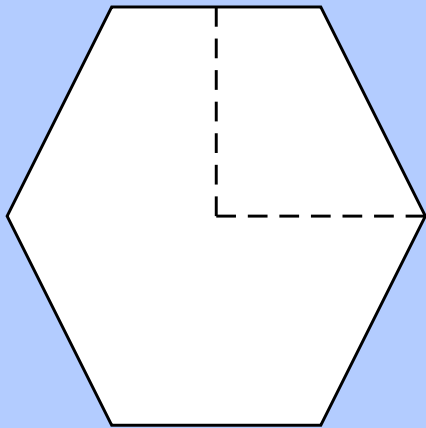
A striped substrate



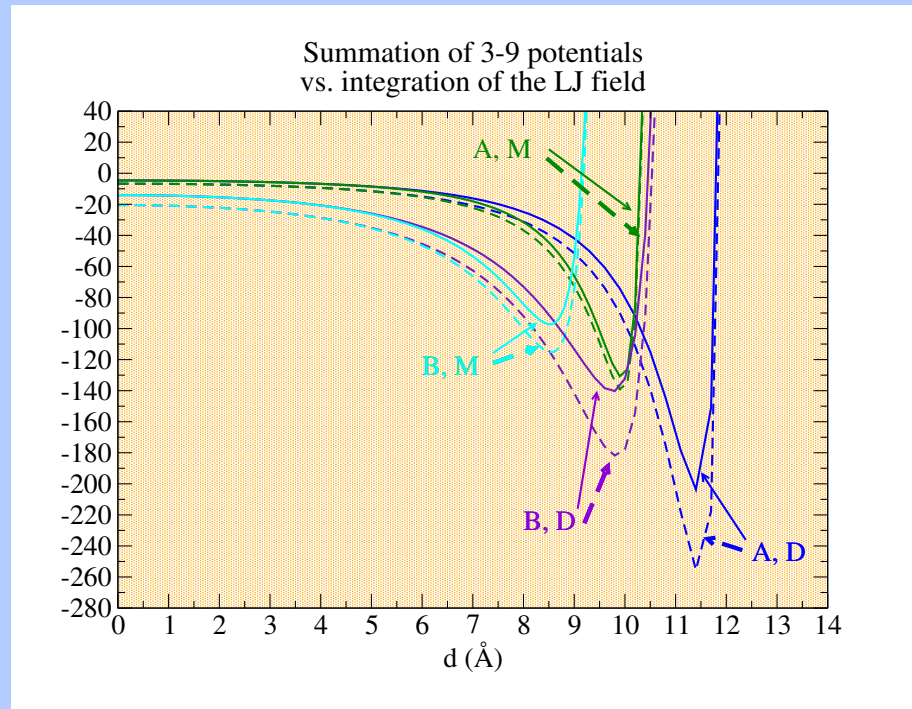
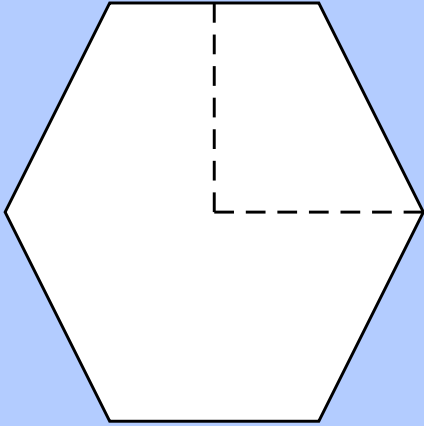
A sawtooth substrate



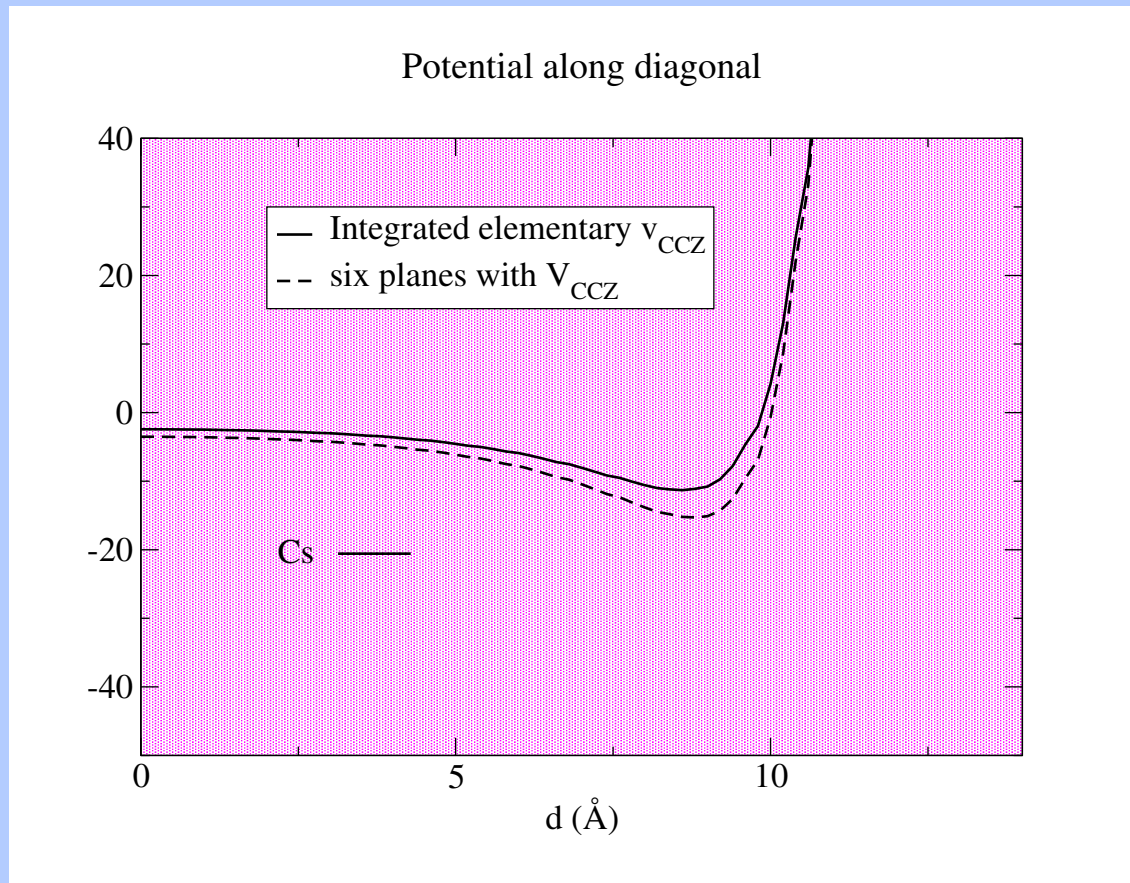
FSM-16 hexagonal pores



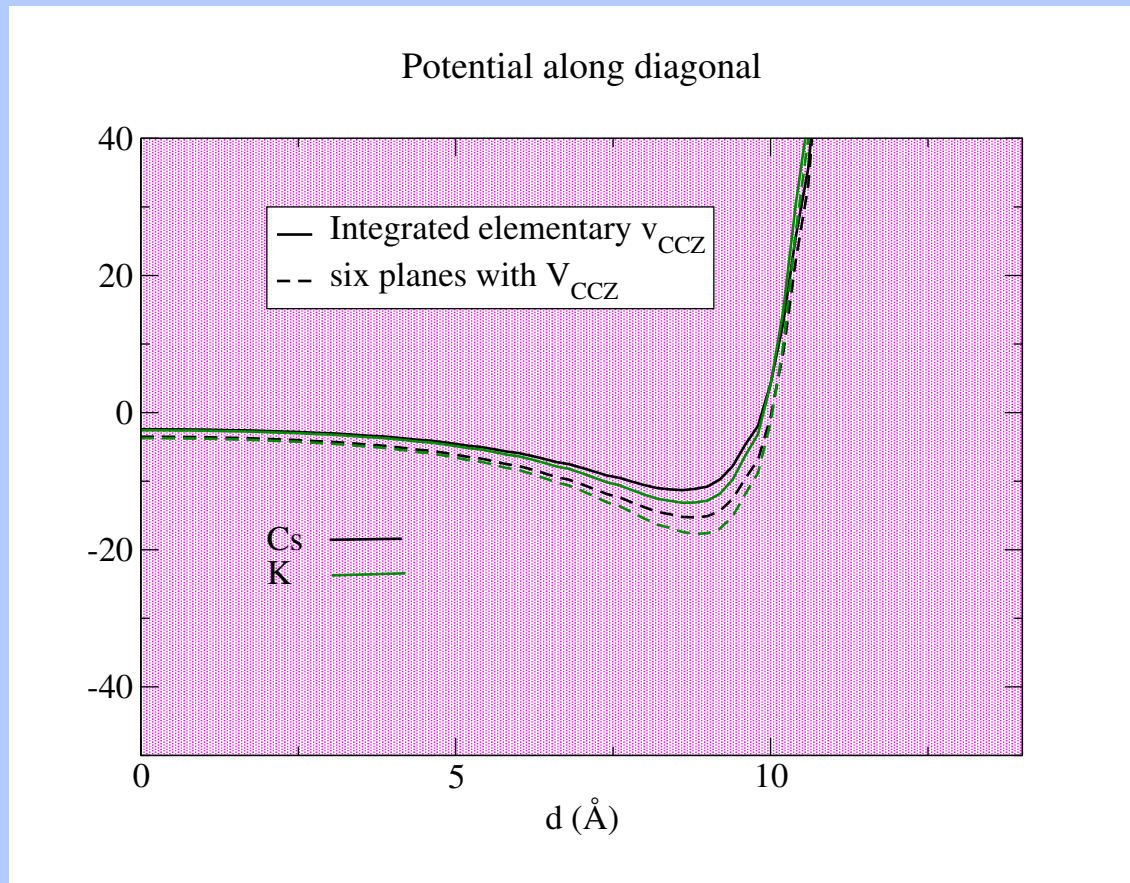
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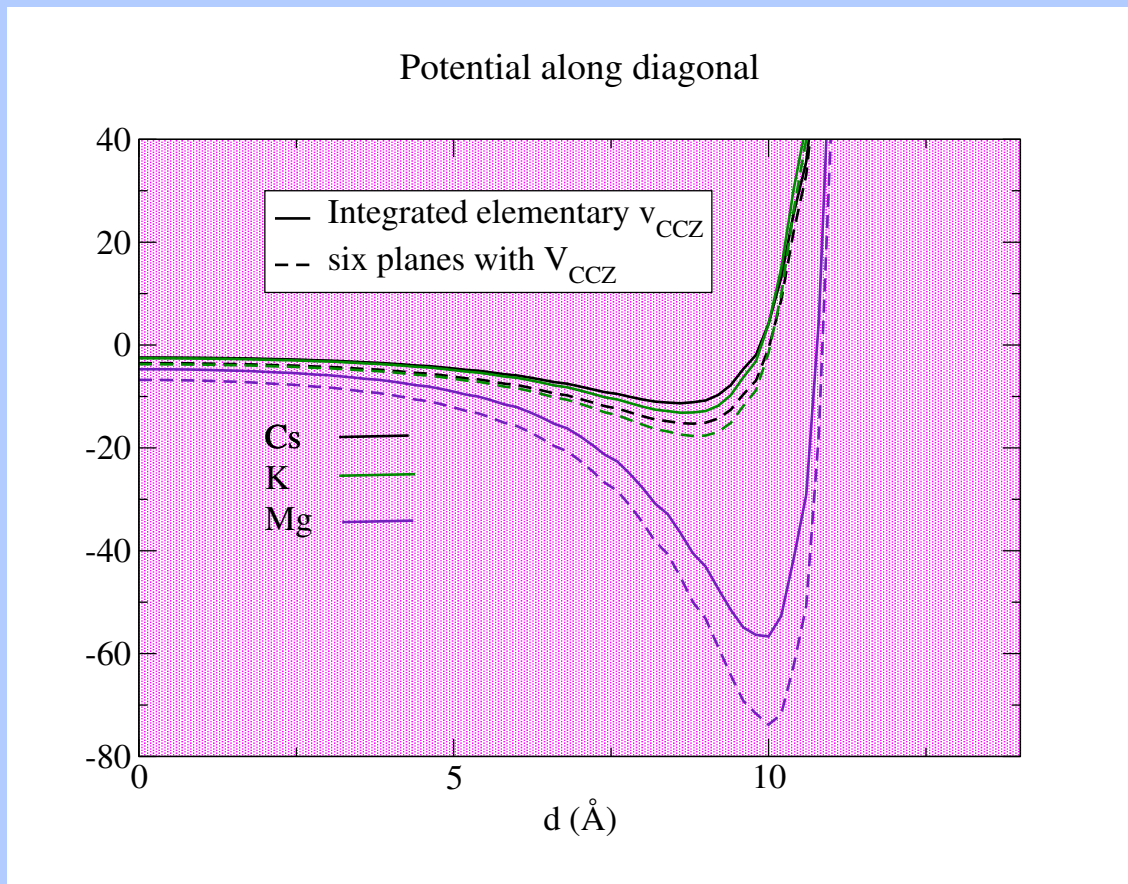
Metallic hexagonal pores



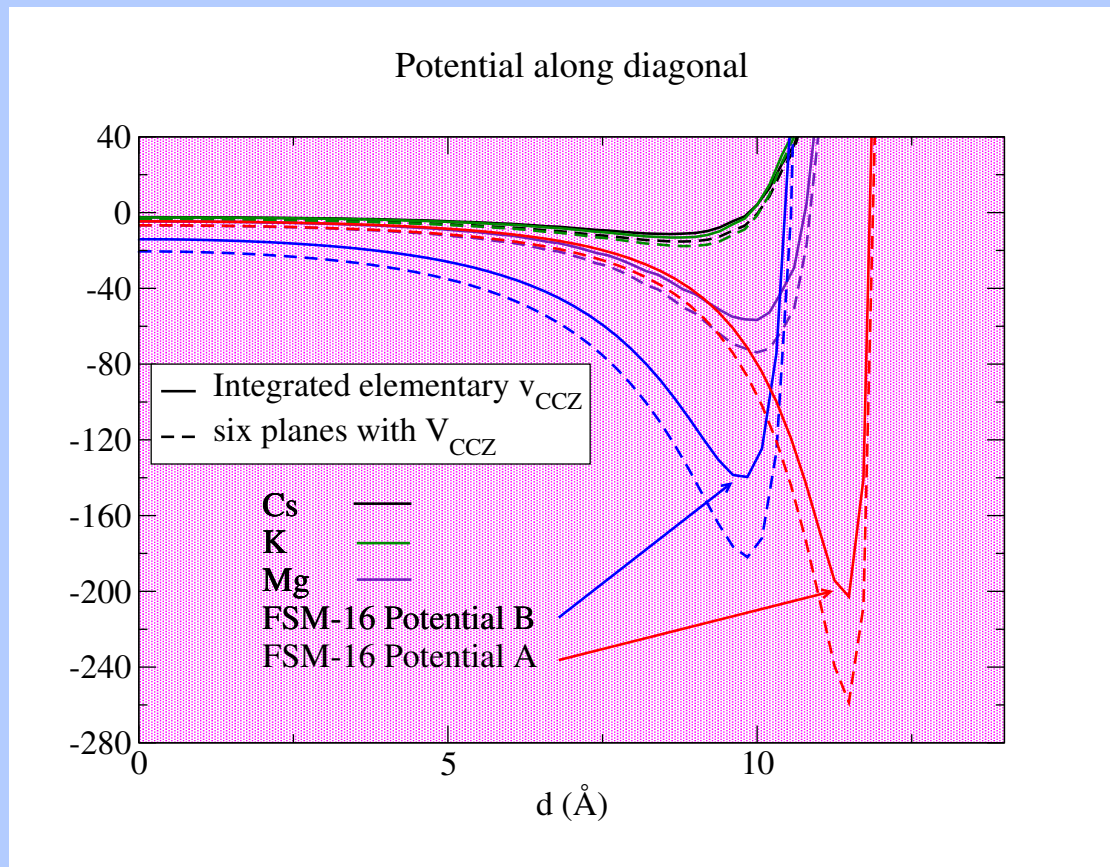
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Standard DF theory at zero temperature

$$\Omega = E - \mu N$$

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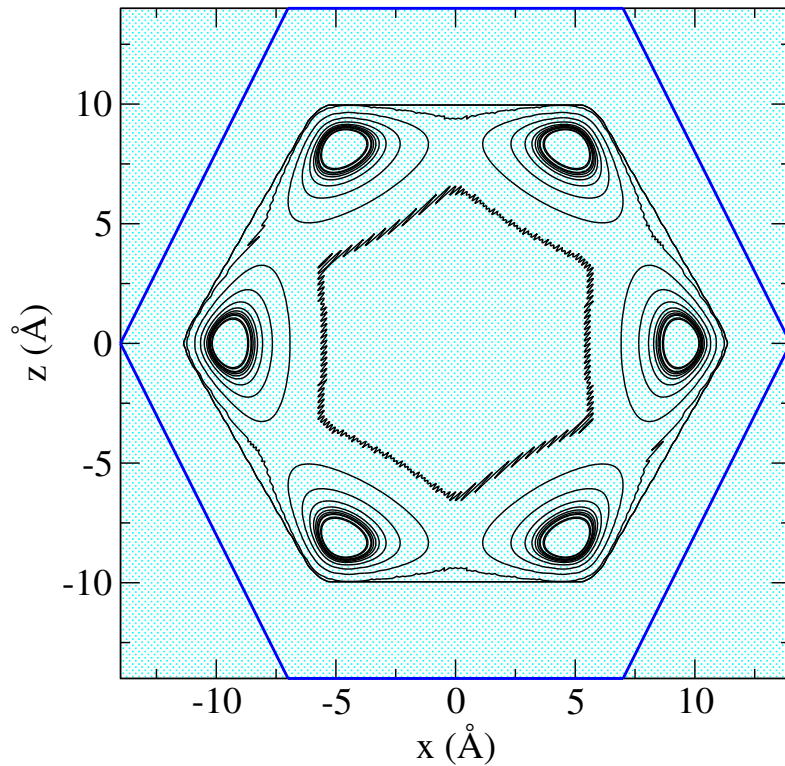
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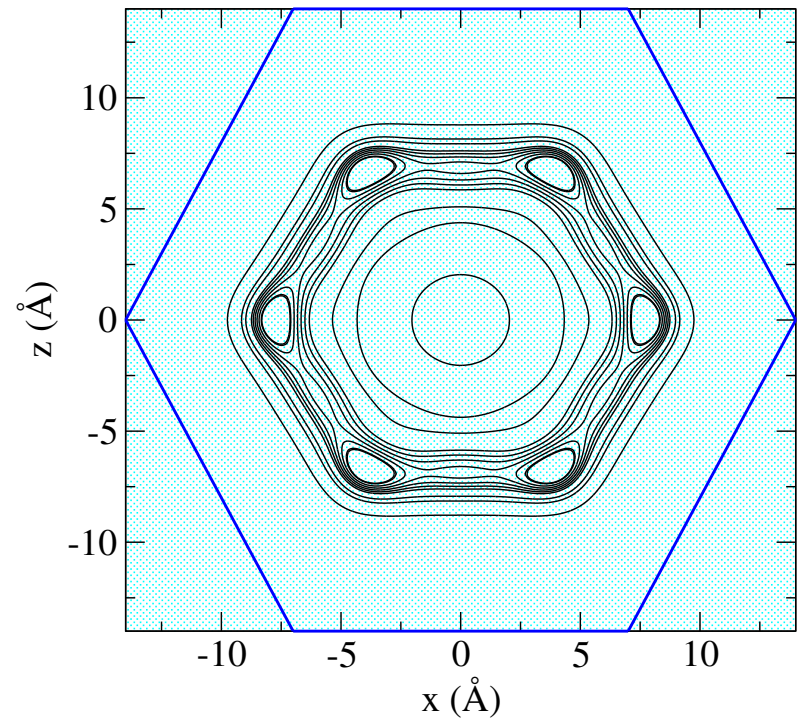
$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(\rho) + V_s(x, z) \right] \sqrt{\rho(x, z)} = \mu \sqrt{\rho(x, z)}$$

Some density profiles obtained with DFT

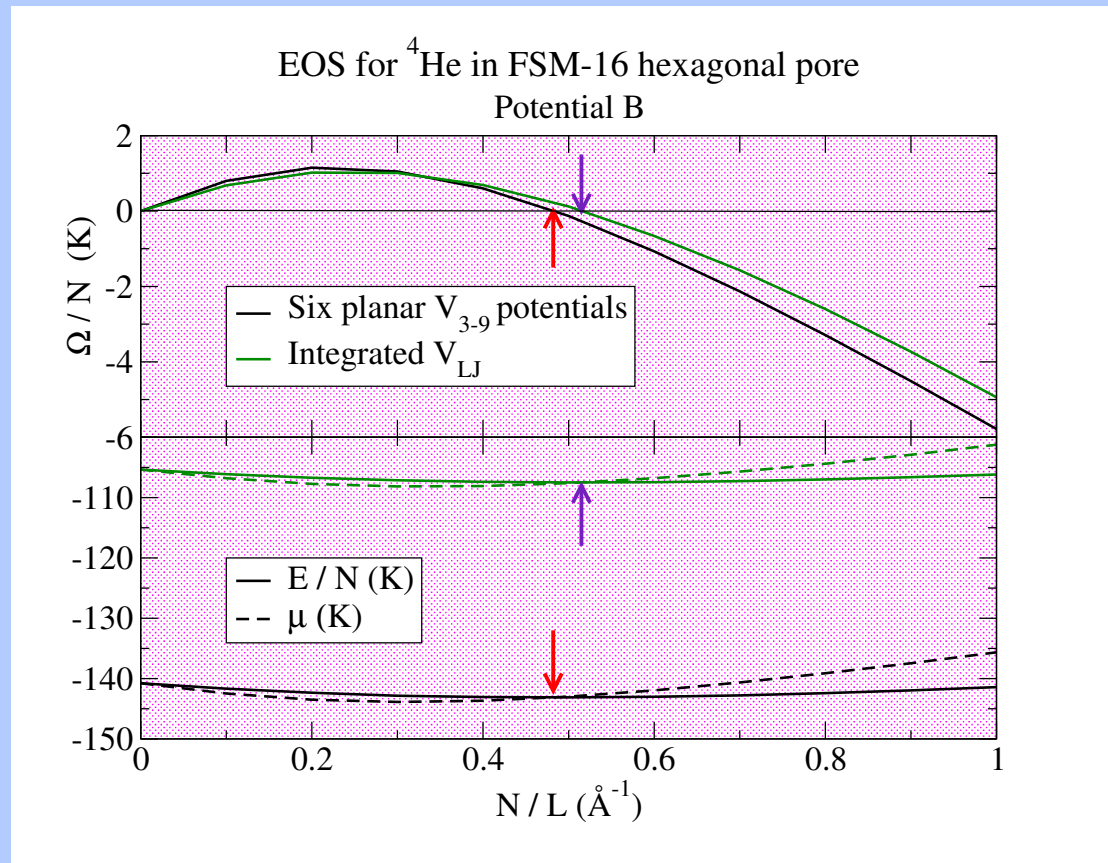
${}^4\text{He}$ in FSM-16 pore
 $N/L = 0.8 \text{ \AA}^{-1}$



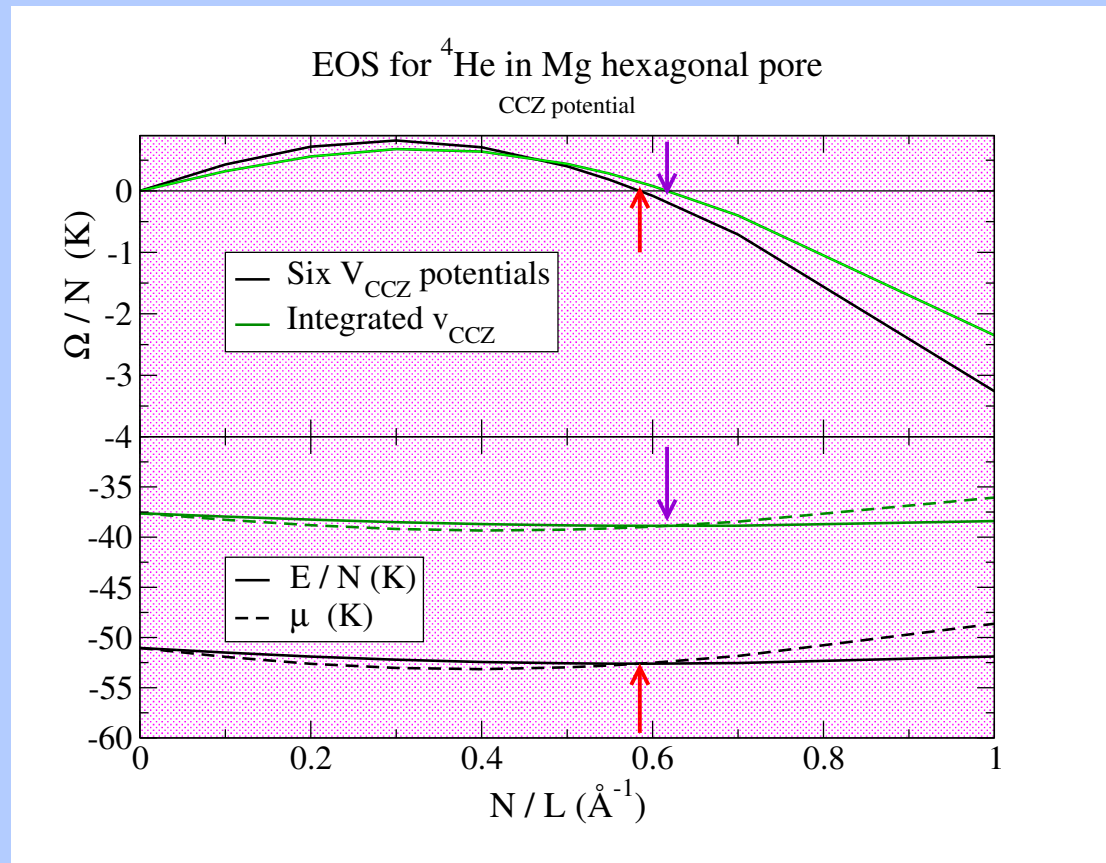
${}^4\text{He}$ in Cs pore
 $N/L = 2 \text{ \AA}^{-1}$



Some EOS computed with DFT



Some EOS computed with DFT



Some energetics

<i>SUBSTRATE</i>	$\varepsilon_{\Sigma}(K)$	$\varepsilon_f(K)$	$\frac{\varepsilon_{\Sigma}-\varepsilon_f}{\varepsilon_f}$
<i>FSM</i> – –16, <i>A</i>	-178.36	-135.98	0.31
<i>FSM</i> – –16, <i>B</i>	-140.77	-105.36	0.34
<i>Mg</i>	-51.04	-37.62	0.36
<i>Li</i>	-23.98	-17.22	0.39
<i>Na</i>	-16.35	-11.58	0.41
<i>K</i>	-10.41	-7.24	0.44
<i>Rb</i>	-9.45	-6.54	0.44
<i>Cs</i>	-9.01	-6.22	0.45

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(J. Klier, P. Leiderer, D. Reinelt, and A. F. G. Wyatt, Phys. Rev. B 72, 245410 (2005))

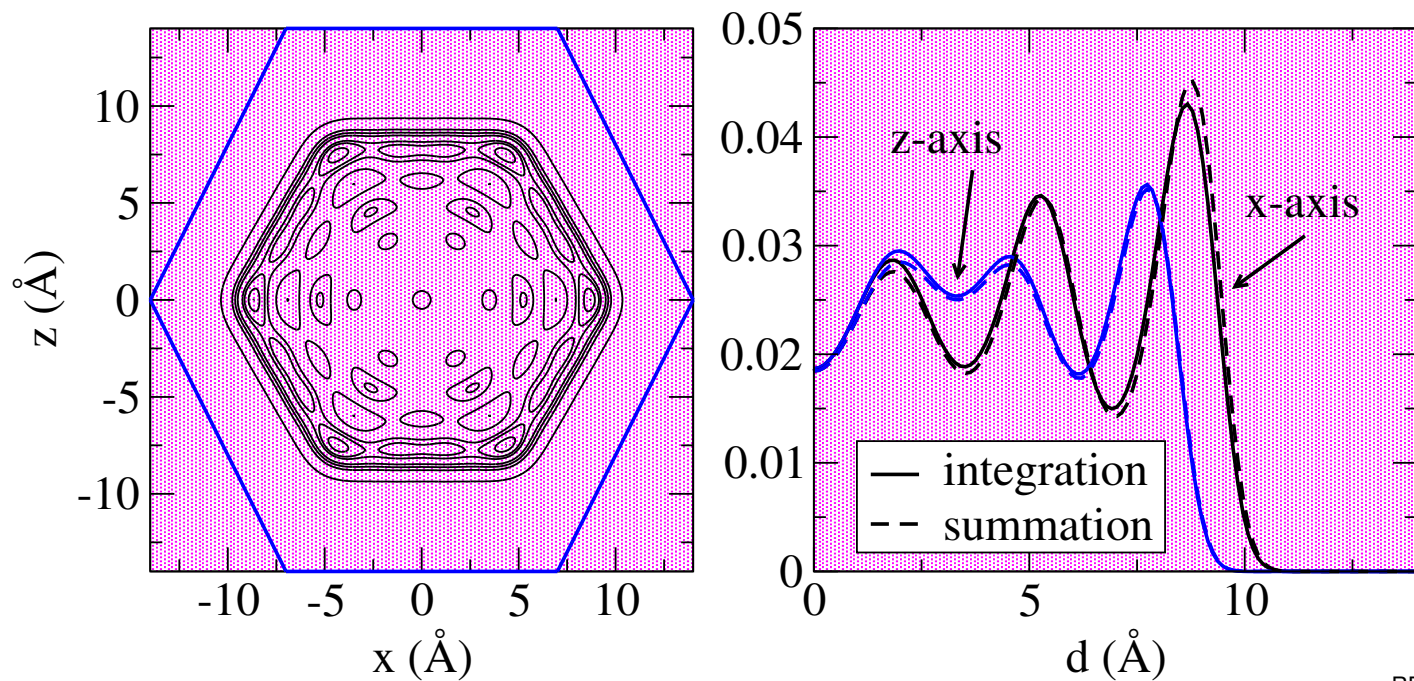
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- \implies the **He–Cs** system is an interesting laboratory for experimenting on wetting physics.

Some comparison among methods

He on Cs Hexagonal pore

at $N/L = 7 \text{ \AA}^{-1}$



Some summary

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- Nextcoming release (RPMBT15 and/or earlier meetings): Condensation of ^4He in polygonal and curved pores in 2D and 3D at zero and finite temperatures.